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Dirac-harmonic maps in dimension 2 with
applications to the Dirac-harmonic heat flow

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BLOW-UP ANALYSIS FOR APPROXIMATE DIRAC-HARMONIC MAPS IN DIMENSION 2 WITH APPLICATIONS TO THE DIRAC-HARMONIC HEAT FLOW

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ABSTRACT. Dirac-harmonic maps couple a second order harmonic map type system with a first nonlinear Dirac equation. We consider approximate Dirac-harmonic maps $\{(\phi_n, \psi_n)\}$, that is, maps that satisfy the Dirac-harmonic system up to controlled error terms. We show that such approximate Dirac-harmonic maps defined on a Riemann surface, that is, in dimension 2, continue to satisfy the basic properties of blow-up analysis like the energy identity and the no neck property. The assumptions are such that they hold for solutions of the heat flow of Dirac-harmonic maps. That flow turns the harmonic map type system into a parabolic system, but simply keeps the Dirac equation as a nonlinear first order constraint along the flow. As a corollary of the main result of this paper, when such a flow blows up at infinite time at interior points, we obtain an energy identity and the no neck property.

1. INTRODUCTION

Dirac-harmonic maps were introduced and studied in [2, 3]. They were motivated by the supersymmetric nonlinear sigma model from quantum field theory [6], and they combine and generalize the theories of harmonic maps and harmonic spinors.

Let us recall the precise definition. Let M be a compact Riemann surface, equipped with a Riemannian metric h and with a fixed spin structure, ΣM be the spinor bundle over M and $\langle \cdot, \cdot \rangle_{\Sigma M}$ be the metric on ΣM induced by the Riemannian metric h . Choosing a local orthonormal basis $e_\alpha, \alpha = 1, 2$ on M , the usual Dirac operator is defined as $\not{D} := e_\alpha \cdot \nabla_{e_\alpha}$, where ∇ is the spin connection on ΣM . The usual Dirac operator \not{D} on a surface can be seen as the Cauchy-Riemann operator. Consider \mathbb{R}^2 with the Euclidean metric $dx^2 + dy^2$. Let $e_1 = \frac{\partial}{\partial x}$ and $e_2 = \frac{\partial}{\partial y}$ be the standard orthonormal frame. A spinor field is simply a map $\psi : \mathbb{R}^2 \rightarrow \Delta_2 = \mathbb{C}^2$, and the action of e_1 and e_2 on spinors can be identified with multiplication with matrices

$$e_1 = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}, \quad e_2 = \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix}.$$

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If $\psi := \begin{pmatrix} \psi_1 \\ \psi_2 \end{pmatrix} : \mathbb{R}^2 \rightarrow \mathbb{C}^2$ is a spinor field, then the Dirac operator is

$$(1.1) \quad \not{D}\psi = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \begin{pmatrix} \frac{\partial \psi_1}{\partial x} \\ \frac{\partial \psi_2}{\partial x} \end{pmatrix} + \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix} \begin{pmatrix} \frac{\partial \psi_1}{\partial y} \\ \frac{\partial \psi_2}{\partial y} \end{pmatrix} = 2 \begin{pmatrix} \frac{\partial \psi_2}{\partial \bar{z}} \\ -\frac{\partial \psi_1}{\partial z} \end{pmatrix},$$

where

$$\frac{\partial}{\partial z} = \frac{1}{2} \left(\frac{\partial}{\partial x} - i \frac{\partial}{\partial y} \right), \quad \frac{\partial}{\partial \bar{z}} = \frac{1}{2} \left(\frac{\partial}{\partial x} + i \frac{\partial}{\partial y} \right).$$

For more details on spin geometry and Dirac operators, one can refer to [13].

Let ϕ be a smooth map from M to another compact Riemannian manifold (N, g) with dimension $n \geq 2$. Let ϕ^*TN be the pull-back bundle of TN by ϕ and then we get the twisted bundle $\Sigma M \otimes \phi^*TN$. Naturally, there is a metric $\langle \cdot, \cdot \rangle_{\Sigma M \otimes \phi^*TN}$ on $\Sigma M \otimes \phi^*TN$ which is induced from the metrics on ΣM and ϕ^*TN . Also we have a natural connection $\tilde{\nabla}$ on $\Sigma M \otimes \phi^*TN$ which is induced from the connections on ΣM and ϕ^*TN . Let ψ be a section of the bundle $\Sigma M \otimes \phi^*TN$. In local coordinates, it can be written as

$$\psi = \psi^i \otimes \partial_{y^i}(\phi),$$

where each ψ^i is a usual spinor on M and ∂_{y^i} is the nature local basis on N . Then $\tilde{\nabla}$ becomes

$$(1.2) \quad \tilde{\nabla}\psi = \nabla\psi^i \otimes \partial_{y^i}(\phi) + (\Gamma_{jk}^i \nabla\phi^j)\psi^k \otimes \partial_{y^i}(\phi),$$

where Γ_{jk}^i are the Christoffel symbols of the Levi-Civita connection of N . The Dirac operator along the map ϕ is defined by $\not{D}\psi := e_\alpha \cdot \tilde{\nabla}_{e_\alpha}\psi$.

We consider the following functional

$$\begin{aligned} L(\phi, \psi) &= \int_M (|d\phi|^2 + \langle \psi, \not{D}\psi \rangle_{\Sigma M \otimes \phi^*TN}) dvol \\ &= \int_M \left(g_{ij}(\phi) h^{\alpha\beta} \frac{\partial \phi^i}{\partial x^\alpha} \frac{\partial \phi^j}{\partial x^\beta} + g_{ij}(\phi) \langle \psi^i, \not{D}\psi^j \rangle_{\Sigma M} \right) dvol. \end{aligned}$$

The functional $L(\phi, \psi)$ is conformally invariant. That is, for any conformal diffeomorphism $f : M \rightarrow M$, setting

$$\tilde{\phi} = \phi \circ f \quad \text{and} \quad \tilde{\psi} = \lambda^{-1/2} \psi \circ f.$$

Then $L(\tilde{\phi}, \tilde{\psi}) = L(\phi, \psi)$. (For the proof, one can refer to [3].) Here λ is the conformal factor of the conformal map f , i.e. $f^*h = \lambda^2 h$. Critical points (ϕ, ψ) are called Dirac-harmonic maps from M to N .

The Euler-Lagrange equations of the functional L are

$$(1.3) \quad (\Delta\phi^i + \Gamma_{jk}^i h^{\alpha\beta} \phi_\alpha^j \phi_\beta^k) \frac{\partial}{\partial y^i}(\phi(x)) = R(\phi, \psi),$$

$$(1.4) \quad \not{D}\psi = 0,$$

where $R(\phi, \psi)$ is defined by

$$R(\phi, \psi) = \frac{1}{2} R_{lij}^m(\phi(x)) \langle \psi^i, \nabla\phi^l \cdot \psi^j \rangle \frac{\partial}{\partial y^m}(\phi(x)).$$

Here R_{lij}^m stands for the Riemann curvature tensor of the target manifold (N, g) . One can refer to [2, 3].

By the Nash embedding theorem, we embed N into \mathbb{R}^K . Then, the critical points (ϕ, ψ) satisfy the Euler-Lagrange equations

$$(1.5) \quad -\Delta\phi = A(\phi)(d\phi, d\phi) - \operatorname{Re}(P(\mathcal{A}(d\phi(e_\alpha), e_\alpha \cdot \psi); \psi)),$$

$$(1.6) \quad \not\partial\psi = \mathcal{A}(d\phi(e_\alpha), e_\alpha \cdot \psi),$$

where $\not\partial$ is the usual Dirac operator, A is the second fundamental form of N in \mathbb{R}^K , and

$$\begin{aligned} \mathcal{A}(d\phi(e_\alpha), e_\alpha \cdot \psi) &:= (\nabla\phi^i \cdot \psi^j) \otimes A(\partial_{y^i}, \partial_{y^j}), \\ \operatorname{Re}(P(\mathcal{A}(d\phi(e_\alpha), e_\alpha \cdot \psi); \psi)) &:= P(A(\partial_{y^l}, \partial_{y^j}); \partial_{y^i}) \operatorname{Re}(\langle \psi^i, d\phi^l \cdot \psi^j \rangle). \end{aligned}$$

Here $P(\xi; \cdot)$ denotes the shape operator, defined by $\langle P(\xi; X), Y \rangle = \langle A(X, Y), \xi \rangle$ for $X, Y \in \Gamma(TN)$, and $\operatorname{Re}(z)$ denotes the real part of $z \in \mathbb{C}$. We refer to [2, 3, 32, 29, 5, 23, 11] for more details.

Denote

$$W^{2,2}(M, N) := \{ \phi \in W^{2,2}(M, \mathbb{R}^K) \text{ with } \phi(x) \in N \text{ for a.e. } x \in M \},$$

$$\begin{aligned} W^{1,4/3}(M, \Sigma M \otimes \phi^*TN) &:= \{ \psi \in W^{1,4/3}(M, \Sigma M \otimes \mathbb{R}^K) \text{ with } \psi(x) \in \Sigma M \otimes \phi^*TN \\ &\text{for a.e. } x \in M \}. \end{aligned}$$

In this paper, with applications for the Dirac-harmonic heat flow in mind, we want to consider pairs (ϕ, ψ) that satisfy the Euler-Lagrange equations up to an error term in L^1 . Here is the precise

Definition 1.1. $(\phi, \psi) \in W^{2,2}(M, N) \times W^{1,4/3}(M, \Sigma M \otimes \phi^*TN)$ is called an approximate Dirac-harmonic map if there exist $(\tau(\phi, \psi), h(\phi, \psi)) \in L^1(M)$ such that

$$(1.7) \quad \tau(\phi, \psi) = \Delta\phi + A(d\phi, d\phi) - \operatorname{Re}(P(\mathcal{A}(d\phi(e_\alpha), e_\alpha \cdot \psi); \psi)),$$

$$(1.8) \quad h(\phi, \psi) = \not\partial\psi - \mathcal{A}(d\phi(e_\alpha), e_\alpha \cdot \psi).$$

Thus, (ϕ, ψ) is a Dirac-harmonic map if and only if $\tau(\phi, \psi) = h(\phi, \psi) = 0$. In the sequel, we shall need to assume that the error terms are in stronger spaces than L^1 , however. See for instance Theorem 1.2.

As for harmonic maps, the conformal invariance of the energy functional L leads to non-compactness of the set of Dirac-harmonic maps in dimension 2. This has been studied extensively by [2, 31, 17], and in [11] for a more general case. For the harmonic map case, we refer to [7, 21, 16, 14, 15, 30, 28]. Roughly speaking, the results of those papers assert that the failure of strong convergence occurs at finitely many concentration points of the energy. At such points, finitely many bubbles (i.e. nonconstant Dirac-harmonic spheres) separate, and the total energies from these bubbles account for the total loss of Dirichlet energies during the process of convergence. Moreover, the image of the remaining the base map and those of the bubbles are connected in the target manifold. This is called the no neck property.

In this paper, we will extend the results from [2, 31, 17] to the approximate Dirac-harmonic maps from a closed Riemann surface M to a compact Riemannian manifold N .

Denote the energy of ϕ on $\Omega \subset M$ by

$$E(\phi; \Omega) = \frac{1}{2} \int_{\Omega} |\nabla \phi|^2 dM,$$

the energy of ψ on $\Omega \subset M$ by

$$E(\psi; \Omega) = \int_{\Omega} |\psi|^4 dM,$$

and the energy of the pair (ϕ, ψ) on $\Omega \subset M$ by

$$E(\psi, \phi; \Omega) = \int_{\Omega} (|\nabla \phi|^2 + |\psi|^4) dM.$$

We shall often omit the domain M from the notation and simply write $E(\phi) = E(\phi; M)$, $E(\psi) = E(\psi; M)$ and $E(\phi, \psi) = E(\phi, \psi; M)$.

Our first main result is

Theorem 1.2. *For a sequence of smooth approximate Dirac-harmonic maps $\{(\phi_n, \psi_n)\}$ from a closed Riemann surface M to a compact Riemannian manifold N with uniformly bounded energy*

$$E(\phi_n, \psi_n) \leq \Lambda < \infty$$

and

$$\|\tau(\phi_n, \psi_n)\|_{L^2} + \|h(\phi_n, \psi_n)\|_{L^4} \leq \Lambda,$$

we assume $(\phi_n, \psi_n) \rightharpoonup (\phi, \psi)$ weakly in $W^{1,2}(M, N) \times L^4(M, \Sigma M \otimes \phi^*TN)$. Defining the blow-up set

$$(1.9) \quad \mathcal{S} := \bigcap_{r>0} \left\{ x \in M \mid \liminf_{n \rightarrow \infty} \int_{D(x,r)} (|d\phi_n|^2 + |\psi_n|^4) \geq \frac{\epsilon_0^2}{2} \right\},$$

where $\epsilon_0 > 0$ is as in Theorem 2.1, then \mathcal{S} is a finite set $\{p_1, \dots, p_I\}$. A subsequence, still denoted by $\{(\phi_k, \psi_k)\}$, converges weakly in $W_{loc}^{2,2}(M \setminus \mathcal{S}) \times W_{loc}^{1,2}(M \setminus \mathcal{S})$ to (ϕ, ψ) and there are a finite set of Dirac-harmonic spheres $(\sigma_i^l, \xi_i^l) : S^2 \rightarrow N$, $i = 1, \dots, I$; $l = 1, \dots, L_i$ such that

$$(1.10) \quad \lim_{n \rightarrow \infty} E(\phi_n) = E(\phi) + \sum_{i=1}^I \sum_{l=1}^{L_i} E(\sigma_i^l),$$

$$(1.11) \quad \lim_{n \rightarrow \infty} E(\psi_n) = E(\psi) + \sum_{i=1}^I \sum_{l=1}^{L_i} E(\xi_i^l),$$

and the image $\phi(M) \cup_{i=1}^I \cup_{l=1}^{L_i} (\sigma_i^l(S^2))$ is a connected set.

As an application of Theorem 1.2, we study the asymptotic behavior at infinite time for the Dirac-harmonic map flow in dimension 2.

For that purpose, we first review the heat flow for Dirac-harmonic maps as introduced and studied in [4, 12] (a different flow has been introduced and studied in [1]). One tries to find $(\phi, \psi) : M \times [0, \infty) \rightarrow N \times \phi^*TN$ that solves

$$(1.12) \quad \begin{cases} \partial_t \phi = \tau(\phi) - \text{Re}(P(\mathcal{A}(d\phi(e_\alpha), e_\alpha \cdot \psi); \psi), & \text{in } M \times (0, \infty)); \\ \not\partial \psi = \mathcal{A}(d\phi(e_\alpha), e_\alpha \cdot \psi), & \text{in } M \times (0, \infty). \end{cases}$$

with the following boundary-initial data:

$$(1.13) \quad \begin{cases} \phi(x, t) = \varphi(x), & \text{on } \partial M \times [0, \infty); \\ \phi(x, 0) = \phi_0(x), & \text{in } M; \\ \mathcal{B}\psi(x, t) = \mathcal{B}\chi(x), & \text{on } \partial M \times [0, \infty); \\ \phi_0(x) = \varphi(x), & \text{on } \partial M, \end{cases}$$

where $\tau(\phi)$ is the tension field of ϕ , M is a compact Riemannian spin manifold with smooth boundary ∂M , $\phi_0 \in W^{1,2}(M, N)$, $\varphi \in C^{2+\alpha}(\partial M; N)$, $\chi \in C^{1+\alpha}(\partial M; \Sigma M \otimes \phi^*TN)$ are given maps and $\mathcal{B} = \mathcal{B}^\pm$ is the chiral boundary operator defined as follows:

$$(1.14) \quad \mathcal{B}^\pm : L^2(M, \Sigma M \otimes \phi^*TN|_{\partial M}) \rightarrow L^2(M, \Sigma M \otimes \phi^*TN|_{\partial M})$$

$$(1.15) \quad \psi \mapsto \frac{1}{2} (Id \pm \vec{n} \cdot G) \cdot \psi,$$

where \vec{n} is the outward unit normal vector field on ∂M , and G is the chiral operator satisfying:

$$(1.16) \quad G^2 = Id, \quad G^* = G, \quad \nabla G = 0, \quad GX \cdot = -X \cdot G,$$

for any $X \in \Gamma(TM)$.

In dimension $\dim M = 2$, [4] established the short-time existence for (1.12) with smooth initial-boundary data (1.13). Later, [12] showed that, under some smallness assumption for $\|\phi_0\|_{H^1} + \|\mathcal{B}\chi\|_{L^2}$, there exists a unique global weak solution to (1.12) with initial-boundary data (1.13), which has at most finitely many singular times and enjoys the property

$$(1.17) \quad E(\phi(t), \psi(t); M) + \int_{M^t} |\partial_t \phi|^2 dM dt \leq C(M, E(\phi_0), \|\mathcal{B}\psi_0\|_{L^2(\partial M)}), \quad \forall 0 \leq t < \infty.$$

It follows from (1.17) that there exists a sequence $t_n \uparrow \infty$ such that $(\phi_n, \psi_n) := (\phi(\cdot, t_n), \psi(\cdot, t_n)) \in W^{2,2}(M, N) \times W^{1, \frac{4}{3}}(M, \Sigma M \otimes \phi^*TN)$ is an approximate Dirac-harmonic map with boundary-data

$$(1.18) \quad \begin{cases} \phi(x) = \varphi(x), & \text{on } \partial M; \\ \mathcal{B}\psi(x) = \mathcal{B}\chi(x), & \text{on } \partial M, \end{cases}$$

which satisfies the assumptions of Theorem 1.2. In fact, $h(\phi_n, \psi_n) = 0$ and

$$\tau(\phi_n, \psi_n) := \partial_t \phi(\cdot, t_n) \text{ satisfying } \|\tau(\phi_n, \psi_n)\|_{L^2} \rightarrow 0.$$

Thus, as an immediate corollary, we obtain

Theorem 1.3. *For $\dim M = 2$ and $\phi_0 \in H^1(M, N)$, $\varphi \in C^{2+\alpha}(\partial M, N)$, $\chi \in C^{1+\alpha}(\partial M, \Sigma M \otimes \varphi^*TN)$, let $(\phi, \psi) : M \times [0, \infty) \rightarrow N \times \phi^*TN$ be a global weak solution of (1.12) and (1.13), which has finitely many singular times and satisfies (1.17). Moreover, we assume that (ϕ, ψ) blows up at infinite time and at interior points. Then there exist $t_n \uparrow \infty$, a Dirac-harmonic map $(\phi_\infty, \psi_\infty) \in C^{2+\alpha}(M, N) \times C^{1+\alpha}(M, \Sigma M \otimes \phi_\infty^*TN)$ with boundary data $\phi_\infty|_{\partial M} = \varphi$ and $\mathcal{B}\psi_\infty|_{\partial M} = \mathcal{B}\chi$, and a nonnegative integer I and finitely many points $\{p_1, \dots, p_I\} \in M$ such that*

$$(1) \quad (\phi_n, \psi_n) := (\phi(\cdot, t_n), \psi(\cdot, t_n)) \rightharpoonup (\phi_\infty, \psi_\infty) \text{ in } W^{1,2}(M, N) \times W^{1, \frac{4}{3}}(M, \Sigma M \times \phi^*TN);$$

$$(2) \quad (\phi_n, \psi_n) \rightarrow (\phi_\infty, \psi_\infty) \text{ in } W_{loc}^{1,2}(M \setminus \{p_1, \dots, p_I\}) \times L_{loc}^4(M \setminus \{p_1, \dots, p_I\});$$

(3) *For $1 \leq i \leq I$, there exist a positive integer L_i and L_i nontrivial Dirac-harmonic spheres $(\sigma_i^l, \xi_i^l) : S^2 \rightarrow N$, $i = 1, \dots, I$; $l = 1, \dots, L_i$ such that*

$$(1.19) \quad \lim_{n \rightarrow \infty} E(\phi_n) = E(\phi_\infty) + \sum_{i=1}^I \sum_{l=1}^{L_i} E(\sigma_i^l),$$

$$(1.20) \quad \lim_{n \rightarrow \infty} E(\psi_n) = E(\psi_\infty) + \sum_{i=1}^I \sum_{l=1}^{L_i} E(\xi_i^l),$$

and

$$(1.21) \quad \lim_{n \rightarrow \infty} \|\phi(\cdot, t_n) - \phi_\infty(\cdot) - \sigma_i^l(\cdot)\|_{L^\infty(M)} = 0.$$

This paper is organized as follows. In Section 2, we shall prove some basic lemmas, called small energy regularity, Pohozaev's identity and removable singularity, so that the expert will readily know what we are talking about, and we shall recall some known results for later use. In Section 3, we shall establish the three circle theorem for approximate Dirac-harmonic maps which ensures the exponential decay of the tangential energy. Our main result Theorem 1.2 will be proved in Section 4.

2. SOME BASIC LEMMAS

In this section, we will prove some basic lemmas and recall some known results which will be used in this paper.

Firstly, we prove a small energy regularity theorem.

Theorem 2.1. *There is a small constant $\epsilon_0 > 0$ such that if $(\phi, \psi) \in W^{2,p}(D, N) \times W^{1,q}(D, \Sigma D \otimes \phi^*TN)$ is an approximate Dirac-harmonic map from the unit disc D in \mathbb{R}^2 to a compact Riemannian manifold (N, g) with $\tau \in L^p$, $h \in L^q$ for some $\frac{4}{3} \leq p \leq 2$, $\frac{8}{5} \leq q \leq 2$, and satisfies*

$$(2.1) \quad E(\phi, \psi; D) = \int_D (|d\phi|^2 + |\psi|^4) dx < \epsilon_0^2,$$

then

$$\begin{aligned}\|\phi - \bar{\phi}\|_{W^{2,p}(D_{\frac{1}{2}})} &\leq C(\|d\phi\|_{L^2(D)} + \|\tau\|_{L^p(D)}), \\ \|\psi\|_{W^{1,q}(D_{\frac{1}{2}})} &\leq C(\|\psi\|_{L^4(D)} + \|h\|_{L^q(D)}),\end{aligned}$$

where $\bar{\phi} := \frac{1}{|D_{1/2}|} \int_{D_{1/2}} \phi dx$ and $C > 0$ is a constant depending only on p, q, Λ, N .

Moreover, by the Sobolev embedding $W^{2,p}(\mathbb{R}^2) \subset C^0(\mathbb{R}^2)$, we have

$$(2.2) \quad \|\phi\|_{Osc(D_{1/2})} = \sup_{x,y \in D_{1/2}} |\phi(x) - \phi(y)| \leq C(\Lambda, N)(\|\nabla\phi\|_{L^2(D)} + \|\tau(u)\|_{L^p(D)}).$$

Proof. Without loss of generality, we assume $\frac{1}{|D_{1/2}|} \int_{D_{1/2}} \phi dx = 0$.

Choosing a cut-off function $\eta \in C_0^\infty(D)$ satisfying $0 \leq \eta \leq 1, \eta|_{D_{3/4}} \equiv 1, |\nabla\eta| + |\nabla^2\eta| \leq C$, by the standard theory of first order elliptic equations, for any $1 < q < 2$, we have

$$\begin{aligned}\|\eta\psi\|_{W^{1,q}(D)} &\leq C\|\not\partial(\eta\psi)\|_{L^q(D)} \\ &\leq C\|\nabla\eta \cdot \psi + \eta\not\partial\psi\|_{L^q(D)} \\ &\leq C(\|\psi\|_{L^q(D)} + \|d\phi\|\|\eta\psi\|_{L^q(D)} + \|h\|_{L^q(D)}) \\ &\leq C\|d\phi\|_{L^2(D)}\|\eta\psi\|_{L^{\frac{2q}{2-q}}(D)} + C(\|\psi\|_{L^q(D)} + \|h\|_{L^q(D)}) \\ &\leq C\epsilon_0\|\eta\psi\|_{L^{\frac{2q}{2-q}}(D)} + C(\|\psi\|_{L^q(D)} + \|h\|_{L^q(D)}).\end{aligned}$$

Taking $q = \frac{8}{5}$ and $\epsilon_0 > 0$ sufficiently small, by Sobolev embedding, we obtain

$$(2.3) \quad \|\eta\psi\|_{L^8(D)} \leq C\|\eta\psi\|_{W^{1,8/5}(D)} \leq C(\|\psi\|_{L^q(D)} + \|h\|_{L^q(D)}).$$

Computing directly, we get

$$\begin{aligned}|\Delta(\eta\phi)| &= |\eta\Delta\phi + 2\nabla\eta\nabla\phi + \phi\Delta\eta| \\ &\leq C(|\phi| + |d\phi| + |d\phi|\|\eta d\phi\| + |\psi|^2\|\eta d\phi\| + |\tau|) \\ (2.4) \quad &\leq C|d\phi|\|d(\eta\phi)\| + C(|\phi| + |d\phi| + \eta|\psi|^2|d\phi| + |\tau|).\end{aligned}$$

By standard elliptic estimates,

$$\begin{aligned}\|\eta\phi\|_{W^{2,4/3}(D)} &\leq C\|d\phi\|_{L^4(D)}\|d(\eta\phi)\|_{L^4(D)} + C(\|d\phi\|_{L^4(D)} + \|\eta\psi\|^2\|d\phi\|_{L^4(D)} + \|\tau\|_{L^4(D)}) \\ &\leq C\|d(\eta\phi)\|_{L^4(D)}\|d\phi\|_{L^2(D)} + C(\|d\phi\|_{L^4(D)} + \|\eta\psi\|_{L^8(D)}^2\|d\phi\|_{L^2(D)} + \|\tau\|_{L^4(D)}) \\ &\leq C\epsilon_0\|d(\eta\phi)\|_{L^4(D)} + C(\|d\phi\|_{L^2(D)} + \|\tau\|_{L^4(D)}).\end{aligned}$$

Taking $\epsilon_0 > 0$ sufficiently small, we have

$$(2.5) \quad \|d\phi\|_{L^4(D_{3/4})} \leq C\|\eta\phi\|_{W^{2,4/3}(D)} \leq C(\|d\phi\|_{L^2(D)} + \|\tau\|_{L^4(D)}).$$

By the interior elliptic estimates, for any $\frac{4}{3} \leq p \leq 2$ we have

$$\begin{aligned}\|\phi\|_{W^{2,p}(D_{1/2})} &\leq C(\|\Delta\phi\|_{L^p(D_{3/4})} + \|d\phi\|_{L^p(D_{3/4})}) \\ &\leq C(\|d\phi\|_{L^4(D_{3/4})}^2 + \|\psi\|_{L^8(D_{3/4})}^2\|d\phi\|_{L^4(D_{3/4})} + \|\tau\|_{L^p(D)} + \|d\phi\|_{L^p(D_{3/4})}) \\ &\leq C(\|d\phi\|_{L^2(D)} + \|\tau\|_{L^p(D)}).\end{aligned}$$

By the interior elliptic estimates for first order equations, for any $\frac{8}{5} \leq q \leq 2$, we get

$$\begin{aligned} \|\psi\|_{W^{1,q}(D_{1/2})} &\leq C(\|\not{D}\psi\|_{L^q(D_{3/4})} + \|\psi\|_{L^q(D_{3/4})}) \\ &\leq C(\|d\phi\|_{L^4(D_{3/4})}\|\psi\|_{L^4(D_{3/4})} + \|\psi\|_{L^2(D)} + \|h\|_{L^q(D)}) \\ &\leq C(1 + \|d\phi\|_{L^2(D)} + \|\tau\|_{L^p(D)})(\|\psi\|_{L^4(D)} + \|h\|_{L^q(D)}) \\ &\leq C(\|\psi\|_{L^4(D)} + \|h\|_{L^q(D)}). \end{aligned}$$

□

Secondly, by a direct computation, we have the following proposition.

Proposition 2.2. *Suppose $\phi \in C^2(M, N)$, $\psi \in C^2(M, \Sigma M \otimes \phi^*TN)$. Let $\{e_\alpha\}_{\alpha=1}^2$ be a unit normal basis of TM and $e_\beta \in \Gamma(TM)$ a section satisfying*

$$[e_\beta, e_\alpha] = 0, \quad \alpha = 1, 2,$$

then

$$(2.6) \quad \langle \psi, \tilde{\nabla}_{e_\beta}(\not{D}\psi) \rangle = 2\langle \text{Re}(P(\mathcal{A}(d\phi(e_\alpha), e_\alpha \cdot \psi); \psi)), \phi_\beta) \rangle + \langle \psi, \not{D}\psi_\beta \rangle,$$

where $[\cdot, \cdot]$ is the Lie bracket, $\phi_\beta = d\phi(e_\beta)$ and $\psi_\beta = \tilde{\nabla}_{e_\beta}\psi$.

Proof. Proposition 2.2 is proved in [3]. For the reader's convenience, we recall it here.

On the one hand, we have

$$\begin{aligned} \langle \text{Re}(P(\mathcal{A}(d\phi(e_\alpha), e_\alpha \cdot \psi); \psi)), \phi_\beta) \rangle &= \langle \frac{1}{2}R_{lij}^m(\phi)\langle \psi^i, \nabla\phi^l \cdot \psi^j \rangle \partial_{y^m}, \phi_\beta^p \partial_{y^p} \rangle \\ &= \frac{1}{2}R_{mlij}\langle \psi^i, \nabla\phi^l \cdot \psi^j \rangle \phi_\beta^m. \end{aligned}$$

On the other hand, by a direct computation, we get

$$\tilde{\nabla}_{e_\alpha}\tilde{\nabla}_{e_\beta}\psi - \tilde{\nabla}_{e_\beta}\tilde{\nabla}_{e_\alpha}\psi = R^{\Sigma M}(e_\alpha, e_\beta)\psi^i \otimes \partial y^i + R_{lij}^m\phi_\alpha^i\phi_\beta^j\psi^l \otimes \partial y^m,$$

where $R^{\Sigma M}$ is the curvature operator of the connection on the spinor bundle ΣM . For this curvature, we have (see [13, 6, 3])

$$e_\alpha \cdot R^{\Sigma M}(e_\alpha, X)\psi^i = \frac{1}{2}\text{Ric}(X), \quad \forall X \in \Gamma(M).$$

Thus we obtain

$$\begin{aligned} \langle \psi, e_\alpha \cdot R^{\Sigma M}(e_\alpha, e_\beta)\psi^i \otimes \partial y^i \rangle &= g_{ij}\langle \psi^j, e_\alpha \cdot R^{\Sigma M}(e_\alpha, e_\beta)\psi^i \rangle \\ &= \frac{1}{2}g_{ij}\langle \psi^j, \text{Ric}(e_\beta) \cdot \psi^i \rangle = 0. \end{aligned}$$

Therefore,

$$\begin{aligned} \langle \psi, \not{D}\psi_\beta \rangle &= \langle \psi, e_\alpha \cdot \tilde{\nabla}_{e_\alpha}\tilde{\nabla}_{e_\beta}\psi \rangle \\ &= \langle \psi, \tilde{\nabla}_{e_\beta}(e_\alpha \cdot \tilde{\nabla}_{e_\alpha}\psi) \rangle + R_{lij}^m\phi_\beta^j\langle \psi, \nabla\phi^i \cdot \psi^l \otimes \partial y^m \rangle \\ &= \langle \psi, \tilde{\nabla}_{e_\beta}(\not{D}\psi) \rangle + R_{mlij}\phi_\beta^j\langle \psi^m, \nabla\phi^i \cdot \psi^l \rangle \\ &= \langle \psi, \tilde{\nabla}_{e_\beta}(\not{D}\psi) \rangle - R_{mlij}\langle \psi^i, \nabla\phi^l \cdot \psi^j \rangle \phi_\beta^m. \end{aligned}$$

Then, the conclusion of the proposition follows immediately. \square

Thirdly, we compute Pohozaev's identity for approximate Dirac-harmonic maps.

Lemma 2.3. *Let $D \subset \mathbb{R}^2$ be the unit disk and (ϕ, ψ) be a smooth approximate Dirac-harmonic map, then for any $0 < t < \frac{1}{2}$, we have*

$$(2.7) \quad \begin{aligned} t \int_{\partial D_t} (|\phi_r|^2 - \frac{1}{2}|\nabla\phi|^2) &= \frac{1}{2} \int_{\partial D_t} \langle \psi, r^{-1} \frac{\partial}{\partial \theta} \cdot \psi_\theta \rangle - \frac{1}{2} \int_{D_t} \langle \psi, \mathbb{D}\psi \rangle dx \\ &\quad - \operatorname{Re} \int_{D_t} \langle \mathbb{D}\psi, r\psi_r \rangle dx + \int_{D_t} r\phi_r \tau dx, \end{aligned}$$

where (r, θ) are polar coordinates in D centered at 0, $\phi_r = d\phi(\frac{\partial}{\partial r})$, $\psi_r = \tilde{\nabla}_{\frac{\partial}{\partial r}} \psi$ and $\psi_\theta = \tilde{\nabla}_{\frac{\partial}{\partial \theta}} \psi$.

Proof. Multiplying the equation (1.7) by $r\phi_r$ and integrating over D_t , by Proposition 2.2 and the fact that $r\phi_r = x^\beta \frac{\partial \phi}{\partial x^\beta} = x^\beta d\phi(\frac{\partial}{\partial x^\beta})$, we get

$$\begin{aligned} \int_{D_t} r\phi_r \tau dx &= \int_{D_t} r\phi_r \Delta \phi dx - \int_{D_t} \langle r\phi_r, \operatorname{Re}(P(\mathcal{A}(d\phi(e_\alpha), e_\alpha \cdot \psi); \psi)) \rangle dx \\ &= \int_{D_t} r\phi_r \Delta \phi dx + \frac{1}{2} \int_{D_t} \langle x^\beta \psi, \mathbb{D}\psi_\beta \rangle dx - \frac{1}{2} \int_{D_t} \langle x^\beta \psi, \tilde{\nabla}_{\frac{\partial}{\partial x^\beta}} \mathbb{D}\psi \rangle dx \\ &:= \text{II} + \text{III} + \text{III}. \end{aligned}$$

On the one hand, integrating by parts, we have

$$\begin{aligned} \text{II} &= \int_{\partial D_t} r|\phi_r|^2 - \int_{D_t} \nabla \phi \nabla (r\phi_r) dx \\ &= \int_{\partial D_t} r|\phi_r|^2 - \int_{D_t} |\nabla \phi|^2 dx - \frac{1}{2} \int_{D_t} r \frac{\partial}{\partial r} |\nabla \phi|^2 dx \\ &= r \int_{\partial D_t} (|\phi_r|^2 - \frac{1}{2}|\nabla \phi|^2). \end{aligned}$$

On the other hand, by Lemma 2.6, we get

$$(2.8) \quad \begin{aligned} 2\text{III} &= \int_{D_t} \langle x^\beta \psi, \mathbb{D}\psi_\beta \rangle dx \\ &= \int_{D_t} \langle \mathbb{D}(x^\beta \psi), \psi_\beta \rangle dx - \int_{\partial D_t} \langle \vec{n} \cdot x^\beta \psi, \psi_\beta \rangle dx \\ &= - \int_{D_t} \langle \psi, \mathbb{D}\psi \rangle dx + \int_{D_t} \langle \mathbb{D}\psi, r\psi_r \rangle dx + \int_{\partial D_t} \langle \psi, r \frac{\partial}{\partial r} \cdot \psi_r \rangle, \end{aligned}$$

and

$$\begin{aligned}
2\text{IIII} &= - \int_{D_t} \langle x^\beta \psi, \tilde{\nabla}_{\partial_\beta} \mathcal{D}\psi \rangle dx \\
&= - \int_{\partial D_t} \langle r\psi, \mathcal{D}\psi \rangle dx + \int_{D_t} \langle \tilde{\nabla}_{\partial_\beta} (x^\beta \psi), \mathcal{D}\psi \rangle dx \\
(2.9) \quad &= 2 \int_{D_t} \langle \psi, \mathcal{D}\psi \rangle dx + \int_{D_t} \langle r\psi_r, \mathcal{D}\psi \rangle dx - \int_{\partial D_t} \langle r\psi, \mathcal{D}\psi \rangle.
\end{aligned}$$

Combining the preceding equations, we get the conclusion of the lemma. \square

Corollary 2.4. *Under the assumption of Lemma 2.3, if $\|\tau(\phi, \psi)\|_{L^2(D)} + \|h(\phi, \psi)\|_{L^4(D)} \leq C$, then for any $0 < t < \frac{1}{4}$, we have*

$$(2.10) \quad \int_{D_{2t} \setminus D_t} (|\phi_r|^2 - \frac{1}{2} |\nabla \phi|^2) dx \leq C (\|r^{-1} \frac{\partial \psi}{\partial \theta}\|_{L^{\frac{4}{3}}(D_{2t} \setminus D_t)} + \|r^{-1} \frac{\partial \phi}{\partial \theta}\|_{L^2(D_{2t} \setminus D_t)}) + Ct.$$

Proof. By Lemma 2.3, for any $0 < s < \frac{1}{2}$, we have

$$(2.11) \quad s \int_{\partial D_s} (|\phi_r|^2 - \frac{1}{2} |\nabla \phi|^2) := \mathbb{J}_1 + \dots + \mathbb{J}_4.$$

It is easy to see that

$$\begin{aligned}
\mathbb{J}_2 &= -\frac{1}{2} \int_{D_s} \langle \psi, \mathcal{D}\psi \rangle dx \leq Cs \|\psi\|_{L^4(D_s)} \|h\|_{L^4(D_s)} \leq Cs, \\
\mathbb{J}_4 &= \int_{D_s} r \phi_r \tau dx \leq Cs \|d\phi\|_{L^2(D_s)} \|\tau\|_{L^2(D_s)} \leq Cs,
\end{aligned}$$

and

$$\begin{aligned}
\mathbb{J}_3 &= -\text{Re} \int_{D_s} \langle \mathcal{D}\psi, r\psi_r \rangle dx \leq Cs \|\psi_r\|_{L^{\frac{4}{3}}(D_s)} \|h\|_{L^4(D_s)} \\
&\leq Cs (\|\nabla \psi\|_{L^{\frac{4}{3}}(D_s)} + \|\psi\|_{L^4(D_s)} \|d\phi\|_{L^2(D_s)}) \|h\|_{L^4(D_s)} \\
&\leq Cs,
\end{aligned}$$

where we used the fact that

$$|\psi_r| \leq C (|\frac{\partial \psi}{\partial r}| + |\psi| |\frac{\partial \phi}{\partial r}|), \quad |\psi_\theta| \leq C (|\frac{\partial \psi}{\partial \theta}| + |\psi| |\frac{\partial \phi}{\partial \theta}|).$$

Multiplying (2.7) by $\frac{1}{s}$ and integrating from t to $2t$, we get

$$\begin{aligned}
\int_{D_{2t} \setminus D_t} (|\phi_r|^2 - \frac{1}{2} |\nabla \phi|^2) &\leq \int_t^{2t} \frac{1}{2s} \int_{\partial D_s} \langle \psi, r^{-1} \frac{\partial}{\partial \theta} \cdot \psi_\theta \rangle d\theta ds + Ct \\
&\leq C \|r^{-1} \psi_\theta\|_{L^{\frac{4}{3}}(D_{2t} \setminus D_t)} \|\psi\|_{L^4(D_{2t} \setminus D_t)} + Ct \\
&\leq C (\|r^{-1} \frac{\partial \psi}{\partial \theta}\|_{L^{\frac{4}{3}}(D_{2t} \setminus D_t)} + \|r^{-1} \frac{\partial \phi}{\partial \theta}\|_{L^2(D_{2t} \setminus D_t)}) + Ct.
\end{aligned}$$

\square

Thirdly, we state an interior removable singularity result.

Theorem 2.5. *Let $(\phi, \psi) \in W_{loc}^{2,2}(D \setminus \{0\}) \times W_{loc}^{1,2}(D \setminus \{0\})$ be an approximate Dirac-harmonic map from $D \setminus \{0\}$ to N with finite energy*

$$\|d\phi\|_{L^2(D)} + \|\psi\|_{L^4(D)} \leq C$$

that satisfies

$$(2.12) \quad \tau = f \in L^2(D), \quad x \in D \setminus \{0\},$$

$$(2.13) \quad h = g \in L^2(D), \quad x \in D \setminus \{0\},$$

then (ϕ, ψ) can be extended to a field in $W^{2,2}(D) \times W^{1,2}(D)$.

Proof. By a standard argument as in Lemma A.2 in [9], it is easy to see that (ϕ, ψ) is a weak solution of (2.12) and (2.13). It is known that the equation of ϕ can be written as an elliptic system with an anti-symmetric potential [29, 5, 23]:

$$\Delta\phi = \Omega \cdot \nabla\phi + f,$$

with $\Omega \in L^2(B_1, so(N) \otimes \mathbb{R}^2)$ satisfying $|\Omega| \leq C(|\nabla\phi| + |\psi|^2)$. Then it follows from Riviere's regularity result and its extensions (see [26, 27, 25]) that $\phi \in W^{2,p}(D)$ for any $1 < p < 2$ which implies $\phi \in W^{1,4}(D)$. Applying a simple argument to the Dirac equation for ψ gives that $\psi \in W^{1,2}(D)$. This indicate that $\psi \in L^8(D)$. Then by (2.12), we have $\Delta\phi \in L^2(D)$ which implies the conclusion of the theorem. \square

In the end of this section, we recall two lemmas which are used in this paper.

Lemma 2.6 ([5]). *For any $\psi, \omega \in W^{1,3/4}(M, \Sigma M \otimes \phi^*TN)$, we have*

$$(2.14) \quad \int_M \langle \psi, \mathcal{D}\omega \rangle = \int_M \langle \mathcal{D}\psi, \omega \rangle - \int_{\partial M} \langle \vec{n} \cdot \psi, \omega \rangle$$

where $\langle \psi, \omega \rangle := h_{ij} \langle \psi^i, \omega^j \rangle$.

Proposition 2.7 ([2]). *Let N be a compact Riemannian manifold. Then there exists a constant $\epsilon_1 = \epsilon_1(N) > 0$ such that if (ϕ, ψ) is a smooth Dirac-harmonic map from the standard sphere S^2 to N satisfying*

$$\int_{S^2} (|d\phi|^2 + |\psi|^4) < \epsilon_1,$$

then both ϕ and ψ are trivial.

3. THREE CIRCLE THEOREM FOR APPROXIMATE DIRAC-HARMONIC MAPS

In this section, we will extend the three circle theorem for Dirac-harmonic maps in [17] to the case of approximate Dirac-harmonic maps. The idea is from Qing-Tian's paper [21], which used a special case of the three circle theorem due to Simon [24] to show that the tangential energy of the sequence in the neck region decays exponentially. The second author in cooperation with H. Yin has extended this idea to some fourth order equations, see [18, 19].

Let us first state the three circle theorem for harmonic functions (see [24, 21, 17]).

Theorem 3.1. *There exists a constant $L > 0$, such that if u is a nontrivial smooth harmonic function defined in $[(i-1)L, (i+2)L] \times S^1$ that satisfies*

$$\int_{\{iL\} \times S^1} u d\theta = \int_{\{(i+1)L\} \times S^1} u d\theta = 0,$$

then

$$(3.1) \quad \|u\|_{L^2([iL, (i+1)L] \times S^1)}^2 < \frac{1}{2} (e^{-L} \|u\|_{L^2([(i-1)L, iL] \times S^1)}^2 + e^{-L} \|u\|_{L^2([(i+1)L, (i+2)L] \times S^1)}^2).$$

Next, we prove an L^2 interior estimate for the following integro-differential equations.

Lemma 3.2. *Suppose $u \in W^{2,2}(D_4 \setminus D_1)$, $v \in W^{1,2}(D_4 \setminus D_1)$ and satisfies*

$$(3.2) \quad \Delta u = A^1 u + A^2 \nabla u + A^3 v + \frac{1}{2\pi} \int_0^{2\pi} A^4 u + A^5 \nabla u + A^6 v d\theta + f_1,$$

$$(3.3) \quad \not\partial v = B^1 u + B^2 \nabla u + B^3 v + \frac{1}{2\pi} \int_0^{2\pi} B^4 u + B^5 \nabla u + B^6 v d\theta + f_2,$$

where

$$(3.4) \quad \sum_{i=1}^6 (\|A_i\|_{L^4(D_4 \setminus D_1)} + \|B_i\|_{L^4(D_4 \setminus D_1)}) \leq \rho \quad \text{and} \quad \sum_{i=1}^2 \|f_i\|_{L^2(D_4 \setminus D_1)} \leq C.$$

Then there exists a positive constant ρ_0 such that if $\rho \leq \rho_0$, there holds

$$(3.5) \quad \begin{aligned} & \|u\|_{W^{2,2}(D_3 \setminus D_2)} + \|v\|_{W^{1,2}(D_3 \setminus D_2)} \\ & \leq C (\|u\|_{L^2(D_4 \setminus D_1)} + \|v\|_{L^2(D_4 \setminus D_1)} + \|f_1\|_{L^2(D_4 \setminus D_1)} + \|f_2\|_{L^2(D_4 \setminus D_1)}). \end{aligned}$$

Proof. The proof is similar to Lemma 3.1 in [17] where $f_1 = f_2 = 0$.

Denote $B_\sigma = D_{3+\sigma} \setminus D_{2-\sigma}$, $0 < \sigma < 1$. Let $\sigma' = \frac{\sigma+1}{2}$. Take a cut-off function $\eta(x) = \eta(|x|)$ with compact support in $B_{\sigma'}$ satisfying $\eta(x) \equiv 1$ in B_σ and $|\nabla \eta| \leq \frac{4}{(1-\sigma)}$ and $|\Delta \eta| \leq \frac{16}{(1-\sigma)^2}$. Computing directly, we get

$$\begin{aligned} \Delta(\eta u) &= \eta \Delta u + 2\nabla \eta \nabla u + \Delta \eta u \\ &= (2\nabla \eta + \eta A^2) \nabla u + (\Delta \eta + \eta A^1) u + \eta A^3 v + \eta f_1 \\ &\quad + \eta \cdot \frac{1}{2\pi} \int_0^{2\pi} A^4 u + A^5 \nabla u + A^6 v d\theta. \end{aligned}$$

By the standard elliptic estimate and Sobolev embedding, we have

$$\begin{aligned}
& \|\eta u\|_{W^{2,2}(D_4)} \\
& \leq C(\|A^1 \eta u\|_{L^2(D_4)} + \|A^2 \eta \nabla u\|_{L^2(D_4)} + \|A^3 \eta v\|_{L^2(D_4)} + \|\eta f_1\|_{L^2(D_4)} \\
& \quad + \|\nabla \eta \nabla u\|_{L^2(D_4)} + \|\Delta \eta u\|_{L^2(D_4)} + \|\eta \cdot \frac{1}{2\pi} \int_0^{2\pi} A^4 u + A^5 \nabla u + A^6 v d\theta\|_{L^2(D_4)}) \\
& \leq C(\|A^1\|_{L^2(D_4 \setminus D_1)} + \|A^4\|_{L^2(D_4 \setminus D_1)} + \|A^2\|_{L^4(D_4 \setminus D_1)} + \|A^5\|_{L^4(D_4 \setminus D_1)}) \|\eta u\|_{W^{2,2}(D_4)} \\
& \quad + C(\|A^3\|_{L^4(D_4 \setminus D_1)} + \|A^6\|_{L^4(D_4 \setminus D_1)}) \|\eta v\|_{W^{1,2}(D_4)} + C(\|A^2 \nabla \eta u\|_{L^2(D_4)} \\
& \quad + \|A^5 \nabla \eta u\|_{L^2(D_4)}) + C(\|\nabla \eta \nabla u\|_{L^2(D_4)} + \|\Delta \eta u\|_{L^2(D_4)} + \|\eta f_1\|_{L^2(D_4)}) \\
& \leq C \rho_0 \|\eta u\|_{W^{2,2}(D_4)} + C \rho_0 \|\eta v\|_{W^{1,2}(D_4)} + C(\|A^2\|_{L^4(D_4 \setminus D_1)} + \|A^5\|_{L^4(D_4 \setminus D_1)}) \|\nabla \eta u\|_{W^{1,2}(D_4)} \\
& \quad + C\left(\frac{\|\nabla u\|_{L^2(B_{\sigma'})}}{1-\sigma} + \frac{\|u\|_{L^2(B_{\sigma'})}}{(1-\sigma)^2} + \|\eta f_1\|_{L^2(D_4)}\right) \\
& \leq C \rho_0 (\|\eta u\|_{W^{2,2}(D_4)} + \|\eta v\|_{W^{1,2}(D_4)}) + C\left(\frac{\|\nabla u\|_{L^2(B_{\sigma'})}}{1-\sigma} + \frac{\|u\|_{L^2(B_{\sigma'})}}{(1-\sigma)^2} + \|\eta f_1\|_{L^2(D_4)}\right).
\end{aligned}$$

Similarly, we can compute

$$\oint(\eta v) = \eta B^1 u + \eta B^2 \nabla u + (\eta B^3 + \nabla \eta) v + \eta \frac{1}{2\pi} \int_0^{2\pi} B^4 u + B^5 \nabla u + B^6 v d\theta + \eta f_2.$$

By the first order elliptic estimate, we have

$$\begin{aligned}
& \|\eta v\|_{W^{1,2}(D_4)} \\
& \leq C(\|B^1 \eta u\|_{L^2(D_4)} + \|B^2 \eta \nabla u\|_{L^2(D_4)} + \|B^3 \eta v\|_{L^2(D_4)} + \|\eta f_2\|_{L^2(D_4)} + \|\nabla \eta v\|_{L^2(D_4)} \\
& \quad + \|\eta \cdot \frac{1}{2\pi} \int_0^{2\pi} B^4 u + B^5 \nabla u + B^6 v d\theta\|_{L^2(D_4)}) \\
& \leq C(\|B^1\|_{L^4(D_4 \setminus D_1)} + \|B^4\|_{L^4(D_4 \setminus D_1)} + \|B^2\|_{L^4(D_4 \setminus D_1)} + \|B^5\|_{L^4(D_4 \setminus D_1)}) \|\eta u\|_{W^{2,2}(D_4)} \\
& \quad + C(\|B^3\|_{L^4(D_4 \setminus D_1)} + \|B^6\|_{L^4(D_4 \setminus D_1)}) \|\eta v\|_{W^{1,2}(D_4)} + C(\|B^2 \nabla \eta u\|_{L^2(D_4)} + \|B^5 \nabla \eta u\|_{L^2(D_4)}) \\
& \quad + C(\|\nabla \eta v\|_{L^2(D_4)} + \|\eta f_2\|_{L^2(D_4)}) \\
& \leq C \rho_0 \|\eta u\|_{W^{2,2}(D_4)} + C \rho_0 \|\eta v\|_{W^{1,2}(D_4)} + C(\|B^2\|_{L^4(D_4 \setminus D_1)} + \|B^5\|_{L^4(D_4 \setminus D_1)}) \|\nabla \eta u\|_{W^{1,2}(D_4)} \\
& \quad + C\left(\frac{\|v\|_{L^2(D_4 \setminus D_1)}}{1-\sigma} + \|\eta f_2\|_{L^2(D_4)}\right) \\
& \leq C \rho_0 (\|\eta u\|_{W^{2,2}(D_4)} + \|\eta v\|_{W^{1,2}(D_4)}) + C\left(\frac{\|\nabla u\|_{L^2(B_{\sigma'})}}{1-\sigma} + \frac{\|u\|_{L^2(B_{\sigma'})}}{(1-\sigma)^2} + \frac{\|v\|_{L^2(D_4 \setminus D_1)}}{1-\sigma}\right. \\
& \quad \left. + \|f_2\|_{L^2(D_4 \setminus D_1)}\right).
\end{aligned}$$

Taking ρ_0 sufficiently small, we get

$$\begin{aligned}
& \|\eta u\|_{W^{2,2}(B_1)} + \|\eta v\|_{W^{1,2}(B_1)} \\
(3.6) \quad & \leq C\left(\frac{\|\nabla u\|_{L^2(B_{\sigma'})}}{1-\sigma} + \frac{\|u\|_{L^2(B_{\sigma'})}}{(1-\sigma)^2} + \frac{\|v\|_{L^2(D_4 \setminus D_1)}}{1-\sigma} + \|f_1\|_{L^2(D_4 \setminus D_1)} + \|f_2\|_{L^2(D_4 \setminus D_1)}\right).
\end{aligned}$$

We now introduce seminorms, and define for $j = 0, 1, 2$

$$\Xi_j = \sup_{0 \leq \sigma \leq 1} (1 - \sigma)^j \|D^j u\|_{L^2(B_\sigma)}.$$

Multiplying (3.6) by $(1 - \sigma)^2$ and noting that $1 - \sigma' = \frac{1 - \sigma}{2}$, we have

$$(3.7) \quad \Xi_2 \leq C (\Xi_1 + \Xi_0 + \|v\|_{L^2(D_4 \setminus D_1)} + \|f_1\|_{L^2(D_4 \setminus D_1)} + \|f_2\|_{L^2(D_4 \setminus D_1)}).$$

We claim now that Ξ_j satisfy an interpolation inequality

$$(3.8) \quad \Xi_1 \leq \epsilon \Xi_2 + \frac{C}{\epsilon} \Xi_0$$

for any $\epsilon > 0$, where $C > 0$ is a universal constant. In fact, by the definition of Ξ_1 , for any $\gamma > 0$, we have

$$\begin{aligned} \Xi_1 &\leq (2 - \sigma_\gamma) \|Du\|_{L^2(B_{\sigma_\gamma})} + \gamma \\ &\leq \epsilon (2 - \sigma_\gamma)^2 \|D^2 u\|_{L^2(B_{\sigma_\gamma})} + \frac{C}{\epsilon} \|u\|_{L^2(B_{\sigma_\gamma})} + \gamma, \end{aligned}$$

where the second inequality is derived from the interpolation Theorem 7.27 (or Theorem 7.28) in [8].

By letting $\gamma \rightarrow 0$, we obtain (3.8). Using (3.8) in (3.7), we then obtain

$$\Xi_2 \leq C (\|u\|_{L^2(D_4 \setminus D_1)} + \|v\|_{L^2(D_4 \setminus D_1)} + \|f_1\|_{L^2(D_4 \setminus D_1)} + \|f_2\|_{L^2(D_4 \setminus D_1)}),$$

this is

$$\|D^2 u\|_{L^2(B_\sigma)} \leq \frac{C}{(1 - \sigma)^2} (\|u\|_{L^2(D_4 \setminus D_1)} + \|v\|_{L^2(D_4 \setminus D_1)} + \|f_1\|_{L^2(D_4 \setminus D_1)} + \|f_2\|_{L^2(D_4 \setminus D_1)}).$$

Taking $\sigma = \frac{1}{2}$, it follows

$$(3.9) \quad \|u\|_{W^{2,2}(B_{1/2})} \leq C (\|u\|_{L^2(D_4 \setminus D_1)} + \|v\|_{L^2(D_4 \setminus D_1)} + \|f_1\|_{L^2(D_4 \setminus D_1)} + \|f_2\|_{L^2(D_4 \setminus D_1)}).$$

Choosing a new cut-off function η in (3.6) and using (3.9), we get

$$(3.10) \quad \|v\|_{W^{1,2}(B_{1/4})} \leq C (\|u\|_{L^2(D_4 \setminus D_1)} + \|v\|_{L^2(D_4 \setminus D_1)} + \|f_1\|_{L^2(D_4 \setminus D_1)} + \|f_2\|_{L^2(D_4 \setminus D_1)}).$$

Then it is easy to see that the lemma follows from (3.9) and (3.10). □

Denote $P_i := D_{e^{(i+1)L}r_2} \setminus D_{e^i L r_2}$ and

$$F_i(u, v) := \int_{P_i} \frac{1}{|x|^2} |u|^2 dx + \int_{P_i} \frac{1}{|x|} |v|^2 dx,$$

where $L > 0$ is the constant in Theorem 3.1.

We have the following three circle theorem:

Theorem 3.3. *Suppose $u \in W^{2,2}(P_{i-1} \cup P_i \cup P_{i+1})$, $v \in W^{1,2}(P_{i-1} \cup P_i \cup P_{i+1})$ satisfy equations (3.2) and (3.3). Then there exists a positive constant ρ_0 , such that if $0 < \rho_1 < \rho_0$ and*

$$(3.11) \quad \max_{i-1, i, i+1} (\| |x| f_1 \|_{L^2(P_j)}^2 + \| |x|^{\frac{1}{2}} f_2 \|_{L^2(P_j)}^2) \leq \rho_1 F_i(u, v),$$

and for any $e^{(i-1)L}r_2 \leq r \leq \frac{1}{2}e^{(i+2)L}r_2$, there hold

$$(3.12) \quad \begin{aligned} & \| |x|^{\frac{3}{2}}(|A^1| + |A^4|) \|_{L^4(D_{2r} \setminus D_r)} + \| |x|(|A^3| + |A^6| + |B^1| + |B^4|) \|_{L^4(D_{2r} \setminus D_r)} \\ & + \| |x|^{\frac{1}{2}}(|A^2| + |A^5| + |B^3| + |B^6|) \|_{L^4(D_{2r} \setminus D_r)} + \| |B^2| + |B^5| \|_{L^4(D_{2r} \setminus D_r)} \leq \rho_1, \end{aligned}$$

and

$$\begin{aligned} & \left| \int_0^{2\pi} u(e^{iL}r_2, \theta) d\theta \right|^2 + \left| \int_0^{2\pi} u(e^{(i+1)L}r_2, \theta) d\theta \right|^2 + \left| \int_0^{2\pi} v(e^{iL}r_2, \theta) d\theta \right|^2 \\ & + \left| \int_0^{2\pi} v(e^{(i+1)L}r_2, \theta) d\theta \right|^2 \leq \rho_1 F_i(u, v) \end{aligned}$$

then, there hold

- (a) $F_{i+1}(u, v) \leq e^{-L}F_i(u, v)$ implies $F_i(u, v) \leq e^{-L}F_{i-1}(u, v)$;
- (b) $F_{i-1}(u, v) \leq e^{-L}F_i(u, v)$ implies $F_i(u, v) \leq e^{-L}F_{i+1}(u, v)$;
- (c) either $F_i(u, v) \leq e^{-L}F_{i-1}(u, v)$ or $F_i(u, v) \leq e^{-L}F_{i+1}(u, v)$.

Proof. Since $F_i(u, v)$, (3.11) and (3.12) are scaling invariant, we may assume $r_2 = 1$ and $i = 2$. If the conclusion does not hold, there exist sequences $\rho_{1k} \rightarrow 0$, A_k^j , B_k^j ($j=1, \dots, 6$), f_{1k} , f_{2k} , u_k , v_k such that u_k and v_k satisfy

$$\begin{aligned} \Delta u_k &= A_k^1 u_k + A_k^2 \nabla u_k + A_k^3 v_k + \frac{1}{2\pi} \int_0^{2\pi} A_k^4 u_k + A_k^5 \nabla u_k + A_k^6 v_k d\theta + f_{1k}, \\ \phi v_k &= B_k^1 u_k + B_k^2 \nabla u_k + B_k^3 v_k + \frac{1}{2\pi} \int_0^{2\pi} B_k^4 u_k + B_k^5 \nabla u_k + B_k^6 v_k d\theta + f_{2k}, \end{aligned}$$

and

$$\begin{aligned} & \left| \int_0^{2\pi} u_k(e^{iL}, \theta) d\theta \right|^2 + \left| \int_0^{2\pi} u_k(e^{(i+1)L}, \theta) d\theta \right|^2 + \left| \int_0^{2\pi} u_k(e^{iL}, \theta) d\theta \right|^2 \\ & + \left| \int_0^{2\pi} u_k(e^{(i+1)L}, \theta) d\theta \right|^2 \leq \rho_{1k} F_i(u_k, v_k). \end{aligned}$$

Moreover, for any $e^{(i-1)L} \leq r \leq e^{(i+2)L}$, A_k^j , B_k^j , f_{1k} , f_{2k} satisfy

$$\begin{aligned} & \| |x|^{\frac{3}{2}}(|A_k^1| + |A_k^4|) \|_{L^4(D_{2r} \setminus D_r)} + \| |x|(|A_k^3| + |A_k^6| + |B_k^1| + |B_k^4|) \|_{L^4(D_{2r} \setminus D_r)} \\ & + \| |x|^{\frac{1}{2}}(|A_k^2| + |A_k^5| + |B_k^3| + |B_k^6|) \|_{L^4(D_{2r} \setminus D_r)} + \| |B_k^2| + |B_k^5| \|_{L^4(D_{2r} \setminus D_r)} \leq \rho_{1k} \end{aligned}$$

and

$$\max_{i-1, i, i+1} (\| |x| f_{1k} \|_{L^2(P_j)}^2 + \| |x|^{\frac{1}{2}} f_{2k} \|_{L^2(P_j)}^2) \leq \rho_{1k} F_i(u, v).$$

But, u_k does not satisfy at least one of the conclusions in (a), (b) and (c).

If (a) does not hold, then we have

$$F_2(u_k, v_k) \geq e^L F_3(u_k, v_k) \quad \text{and} \quad F_2(u_k, v_k) \geq e^{-L} F_1(u_k, v_k);$$

If (b) does not hold, then we have

$$F_2(u_k, v_k) \geq e^L F_1(u_k, v_k) \quad \text{and} \quad F_2(u_k, v_k) \geq e^{-L} F_3(u_k, v_k);$$

If (c) does not hold, then we have

$$F_2(u_k, v_k) \geq e^{-L} \max\{F_1(u_k, v_k), F_3(u_k, v_k)\};$$

In all those three cases, we may get the same conclusion that

$$(3.13) \quad 2F_2(u_k, v_k) \geq e^{-L}(F_1(u_k, v_k) + F_3(u_k, v_k)).$$

Without loss of generality, we assume $F_2(u_k, v_k) = 1$ (if not, we consider $\tilde{u}_k = \frac{u_k}{F_2(u_k, v_k)}$ and $\tilde{v}_k = \frac{v_k}{F_2(u_k, v_k)}$). Then we obtain

$$\|u_k\|_{L^2(P_1 \cup P_2 \cup P_3)} + \|v_k\|_{L^2(P_1 \cup P_2 \cup P_3)} \leq C.$$

By Lemma 3.2, we have $\|u_k\|_{W^{2,2}(P_2)} + \|v_k\|_{W^{1,2}(P_2)} \leq C$. So, there exists a subsequence of (u_k, v_k) (we still denote it by (u_k, v_k)), such that

$$\begin{aligned} u_k &\rightharpoonup u, & v_k &\rightharpoonup v \text{ weakly in } L^2(P_1 \cup P_2 \cup P_3); \\ u_k &\rightarrow u, & v_k &\rightarrow v \text{ strongly in } L^2(P_2). \end{aligned}$$

It is easy to see that u is a harmonic function and v is a holomorphic function in $D_{e^{(i+2)L}} \setminus D_{e^{(i-1)L}}$ and they satisfy

$$\int_{\partial D_{e^{iL}}} u = \int_{\partial D_{e^{(i+1)L}}} u = \int_{\partial D_{e^{iL}}} v = \int_{\partial D_{e^{(i+1)L}}} v = 0.$$

Let $f : \mathbb{R}^1 \times \mathbb{S}^1 \rightarrow \mathbb{R}^2$, $f(t, \theta) = (e^t, \theta)$ $(t, \theta) \in \mathbb{R}^1 \times \mathbb{S}^1$ where $\mathbb{R}^1 \times \mathbb{S}^1$ is equipped with the metric $g = dt^2 + d\theta^2$, which is conformal to the standard Euclidean metric ds^2 on \mathbb{R}^2 . In fact,

$$(f^{-1})^*g = \frac{1}{r^2}ds^2.$$

Then we know that $u \circ f$ is a harmonic function and $e^{\frac{t}{2}}v \circ f$ is a holomorphic function in $[L, 4L] \times S^1$.

According to Theorem 3.1, we know

$$\|u \circ f\|_{L^2([2L, 3L] \times S^1)}^2 < \frac{1}{2}(e^{-L}\|u \circ f\|_{L^2([L, 2L] \times S^1)}^2 + e^{-L}\|u \circ f\|_{L^2([3L, 4L] \times S^1)}^2)$$

and

$$\|e^{\frac{t}{2}}v \circ f\|_{L^2([2L, 3L] \times S^1)}^2 < \frac{1}{2}(e^{-L}\|e^{\frac{t}{2}}v \circ f\|_{L^2([L, 2L] \times S^1)}^2 + e^{-L}\|e^{\frac{t}{2}}v \circ f\|_{L^2([3L, 4L] \times S^1)}^2)$$

which implies

$$\left\| \frac{1}{|x|}u \right\|_{L^2(P_2)}^2 < \frac{1}{2}(e^{-L}\left\| \frac{1}{|x|}u \right\|_{L^2(P_1)}^2 + e^{-L}\left\| \frac{1}{|x|}u \right\|_{L^2(P_3)}^2)$$

and

$$\left\| \frac{1}{\sqrt{|x|}}v \right\|_{L^2(P_2)}^2 < \frac{1}{2}(e^{-L}\left\| \frac{1}{\sqrt{|x|}}v \right\|_{L^2(P_1)}^2 + e^{-L}\left\| \frac{1}{\sqrt{|x|}}v \right\|_{L^2(P_3)}^2).$$

Thus,

$$(3.14) \quad 2F_2(u, v) < e^{-L}(F_1(u, v) + F_3(u, v)).$$

But, letting $k \rightarrow \infty$ in (3.13) which implies

$$(3.15) \quad 2F_2(u, v) \geq e^{-L}(F_1(u, v) + F_3(u, v)).$$

This is a contradiction. This proves the theorem. \square

As a direct corollary of the three circle theorem, we can get the following decay lemma.

Lemma 3.4. *Let $\rho_1 > 0$ be the constant in Theorem 3.3. Let $u \in W^{2,2}(D_{e^{(l+1)L}r_2} \setminus D_{r_2})$, $v \in W^{1,2}(D_{e^{(l+1)L}r_2} \setminus D_{r_2})$, $f_i \in L^2(D_{e^{(l+1)L}r_2} \setminus D_{r_2})$, $i = 1, 2$, and some integer $l > 1$, satisfying equations (3.2), (3.3) and for any $r_2 \leq r \leq \frac{1}{2}e^{(l+1)L}r_2$, there hold*

$$(3.16) \quad \begin{aligned} & \| |x|^{\frac{3}{2}}(|A^1| + |A^4|) \|_{L^4(D_{2r} \setminus D_r)} + \| |x|(|A^3| + |A^6| + |B^1| + |B^4|) \|_{L^4(D_{2r} \setminus D_r)} \\ & + \| |x|^{\frac{1}{2}}(|A^2| + |A^5| + |B^3| + |B^6|) \|_{L^4(D_{2r} \setminus D_r)} + \| |B^2| + |B^5| \|_{L^4(D_{2r} \setminus D_r)} \leq \rho_1, \end{aligned}$$

and

$$\int_{\partial D_r} u = \int_{\partial D_r} v = 0.$$

Then we have

$$(3.17) \quad F_i(u, v) \leq C(e^{Lr_2} + F_0(u, v) + F_l(u, v))(e^{-(l-i)L} + e^{-iL}).$$

Proof. Denote the set of j ($0 < j < l$) for which

$$(3.18) \quad \max_{j-1, j, j+1} (\| |x|f_1 \|_{L^2(P_i)}^2 + \| |x|^{\frac{1}{2}}f_2 \|_{L^2(P_i)}^2) > \rho_1 F_j(u, v)$$

by $J := \{j_1, \dots, j_k\}$. If $J = \emptyset$, according to (c) of Theorem 3.3, we get

$$F_i(u, v) \leq e^{-L}F_{i-1}(u, v) \text{ or } F_i(u, v) \leq e^{-L}F_{i+1}(u, v).$$

Then using the (a) and (b) of Theorem 3.3, by iterating, we obtain

$$F_i(u, v) \leq e^{-Li}F_0(u, v) \text{ or } F_i(u, v) \leq e^{-L(l-i)}F_l(u, v).$$

So, we have

$$(3.19) \quad F_i(u, v) \leq \frac{1}{2}(e^{-Li}F_0(u, v) + e^{-L(l-i)}F_l(u, v)),$$

which implies (3.17) immediately.

If $J \neq \emptyset$, without loss of generality, we may assume

$$0 < j_1 < j_2 < \dots < j_k < l.$$

Then for each j_m , $m = 1, \dots, k$, we have

$$\begin{aligned} F_{j_m}(u, v) & \leq C \max_{j_m-1, j_m, j_m+1} (\| |x|f_1 \|_{L^2(P_i)}^2 + \| |x|^{\frac{1}{2}}f_2 \|_{L^2(P_i)}^2) \\ & \leq Ce^{j_m L}r_2 = C(e^{Lr_2})e^{-L(l-j_m)}. \end{aligned}$$

By the choice of j_m , the condition (3.11) holds for $j_m < i < j_{m+1}$, $m = 1, \dots, k-1$. Similar to deriving (3.19), we obtain

$$F_i(u, v) \leq e^{-L(i-j_m)}F_{j_m}(u, v) \text{ or } F_i(u, v) \leq e^{-L(j_{m+1}-i)}F_{j_{m+1}}(u, v).$$

Thus, we know for $j_1 \leq i \leq j_k$, there exists $m \in \{1, \dots, k-1\}$ such that $j_m \leq i \leq j_{m+1}$. Then we get

$$\begin{aligned}
(3.20) \quad F_i(u, v) &\leq \frac{1}{2}(e^{-L(i-j_m)}F_{j_m}(u, v) + e^{-L(j_{m+1}-i)}F_{j_{m+1}}(u, v)) \\
&\leq C(e^{lL}r_2)(e^{-L(i-j_m)}e^{-L(l-j_m)} + e^{-L(j_{m+1}-i)}e^{-L(l-j_{m+1})}) \\
&\leq C(e^{lL}r_2)e^{-(l-i)L}.
\end{aligned}$$

So, if $j_1 = 1$ and $j_k = l-1$, then the inequality (3.19) follows immediately. If not, assuming $j_1 > 1$, similar to deriving (3.20), we have, for $1 \leq i \leq j_1$,

$$\begin{aligned}
F_i(u, v) &\leq \frac{1}{2}(e^{-Li}F_0(u, v) + e^{-L(j_1-i)}F_{j_1}(u, v)) \\
&\leq C(e^{-Li}F_0(u, v) + (e^{lL}r_2)e^{-(l-i)L}).
\end{aligned}$$

Similarly, if $j_k < l-1$, then for $j_k \leq i \leq l$, we have

$$\begin{aligned}
F_i(u, v) &\leq \frac{1}{2}(e^{-L(l-i)}F_l(u, v) + e^{-L(i-j_k)}F_{j_k}(u, v)) \\
&\leq C(e^{-L(l-i)}F_l(u, v) + (e^{lL}r_2)e^{-(l-i)L}).
\end{aligned}$$

Combining the preceding estimates proves the lemma. \square

Corollary 3.5. *Under the assumptions of Lemma 3.4, we have*

$$\begin{aligned}
(3.21) \quad &\|\nabla u\|_{L^2(P_i)} + \|\nabla v\|_{L^{\frac{4}{3}}(P_i)} \\
&\leq C((e^{lL}r_2)^{\frac{1}{2}} + F_0^{1/2}(u, v) + F_l^{1/2}(u, v))(e^{-\frac{1}{2}(l-i)L} + e^{-\frac{1}{2}iL}).
\end{aligned}$$

Proof. By Lemma 3.2, Lemma 3.4 and a standard scaling argument, we get

$$\begin{aligned}
&\|\nabla u\|_{L^2(P_i)} + \|\nabla v\|_{L^{\frac{4}{3}}(P_i)} \\
&\leq C(F_{i-1}^{1/2}(u, v) + F_i^{1/2}(u, v) + F_{i+1}^{1/2}(u, v) + e^{iL}r_2\|f_1\|_{L^2(P_{i-1} \cup P_i \cup P_{i+1})} \\
&\quad + (e^{iL}r_2)^{\frac{1}{2}}\|f_2\|_{L^2(P_{i-1} \cup P_i \cup P_{i+1})}) \\
&\leq C((e^{lL}r_2)^{\frac{1}{2}} + F_0^{1/2}(u) + F_l^{1/2}(u))(e^{-\frac{1}{2}(l-i)L} + e^{-\frac{1}{2}iL}).
\end{aligned}$$

\square

4. ENERGY IDENTITY AND NO NECK RESULT

In this section, we will prove our main result Theorem 1.2.

We first consider the following simpler case of a single interior blow-up point.

Theorem 4.1. *Let (ϕ_n, ψ_n) be a sequence of smooth approximate Dirac-harmonic maps from $D_1(0)$ to N with*

$$(a) \quad \|\phi_n\|_{W^{1,2}(D)} + \|\psi_n\|_{W^{1,\frac{4}{3}}(D)} + \|\tau_n\|_{L^2(D)} + \|h_n\|_{L^4(D)} \leq \Lambda,$$

(b) $(\phi_n, \psi_n) \rightarrow (\phi, \psi)$ strongly in $W_{loc}^{1,2}(D \setminus \{0\}, \mathbb{R}^N) \times W_{loc}^{1,\frac{4}{3}}(D \setminus \{0\}, \mathbb{R}^N)$ as $n \rightarrow \infty$, where $\tau_n := \tau(\phi_n, \psi_n)$ and $h_n := h(\phi_n, \psi_n)$. Then there exist a subsequence of (ϕ_n, ψ_n) (still denoted by (ϕ_n, ψ_n)) and a nonnegative integer Q such that, for any $i = 1, \dots, Q$, there exist point x_n^i , positive numbers λ_n^i and a nonconstant Dirac-harmonic sphere $(\sigma^i, \xi^i) : S^2 \rightarrow N$ such that:

$$(1) \quad x_n^i \rightarrow 0, \lambda_n^i \rightarrow 0, \text{ as } n \rightarrow \infty;$$

$$(2) \quad \lim_{n \rightarrow \infty} \left(\frac{\lambda_n^i}{\lambda_n^j} + \frac{\lambda_n^j}{\lambda_n^i} + \frac{|x_n^i - x_n^j|}{\lambda_n^i + \lambda_n^j} \right) = \infty \text{ for any } i \neq j;$$

$$(3) \quad (\sigma^i, \xi^i) \text{ is the weak limit of } (\phi_n(x_n^i + \lambda_n^i x), \psi_n(x_n^i + \lambda_n^i x)) \text{ in } W_{loc}^{1,2}(\mathbb{R}^2) \times W_{loc}^{1,\frac{4}{3}}(\mathbb{R}^2);$$

(4) **Energy identity:** we have

$$(4.1) \quad \lim_{n \rightarrow \infty} E(\phi_n) = E(\phi) + \sum_{i=1}^Q E(\sigma^i),$$

$$(4.2) \quad \lim_{n \rightarrow \infty} E(\psi_n) = E(\psi) + \sum_{i=1}^Q E(\xi^i),$$

(5) **No neck property:** The image

$$(4.3) \quad \phi(D) \cup \bigcup_{i=1}^Q \sigma^i(S^2)$$

is a connected set.

Proof. Assume 0 is the only blow-up point of the sequence $\{(\phi_n, \psi_n)\}$ in D , i.e.

$$(4.4) \quad \liminf_{n \rightarrow \infty} E(\phi_n, \psi_n; D_r) \geq \frac{\epsilon_0^2}{2} \text{ for all } r > 0.$$

By the standard argument of blow-up analysis we can assume that, for any n , there exist sequences $x_n \rightarrow 0$ and $r_n \rightarrow 0$ such that

$$(4.5) \quad E(\phi_n, \psi_n; D_{r_n}(x_n)) = \sup_{\substack{x \in D, r \leq r_n \\ D_r(x) \subset D}} E(\phi_n, \psi_n; D_r(x)) = \frac{\epsilon_0^2}{4}.$$

Denoting

$$(4.6) \quad \phi'_n(x) := \phi_n(x_n + r_n x), \psi'_n(x) := \sqrt{r_n} \psi_n(x_n + r_n x)$$

then we have

$$(4.7) \quad \tau(\phi'_n, \psi'_n) = \Delta \phi'_n + A(d\phi'_n, d\phi'_n) - \text{Re}(P(\mathcal{A}(d\phi'_n(e_\alpha), e_\alpha \cdot \psi'_n); \psi'_n)),$$

$$(4.8) \quad h(\phi'_n, \psi'_n) = \not\partial \psi'_n - \mathcal{A}(d\phi'_n(e_\alpha), e_\alpha \cdot \psi'_n),$$

where $\tau(\phi'_n, \psi'_n) = r_n^2 \tau(\phi_n, \psi_n)$ and $h(\phi'_n, \psi'_n) = r_n^{3/2} h(\phi_n, \psi_n)$. Noting that for any $D(y) \subset \mathbb{R}^2$, there holds

$$E(\phi'_n, \psi'_n; D(y)) \leq E(\phi_n, \psi_n; D_{r_n}(x_n)) = \frac{\epsilon_0^2}{4}.$$

By the small energy regularity Theorem 2.1, the removable singularity Theorem 2.5 and conformal invariance of Dirac-harmonic maps in dimension two, we can take a subsequence, still denoted by (ϕ'_n, ψ'_n) , that strongly converges to a nonconstant Dirac-harmonic sphere. This is the first bubble.

By the standard induction argument in [7], we only need to prove the theorem in the case where there is only one bubble. Under this assumption, we have the following:

Claim: for any $\epsilon > 0$, there exist $\delta > 0$ and $R > 0$ such that

$$(4.9) \quad E(\phi_n, \psi_n; D_{8t}(x_n) \setminus D_t(x_n)) \leq \epsilon^2 \text{ for any } t \in (\frac{1}{2}r_n R, 2\delta)$$

when n is large enough.

In fact, if (4.9) is not true, then we can find $\bar{\epsilon} > 0$, $t_n \rightarrow 0$, such that $\lim_{n \rightarrow \infty} \frac{t_n}{r_n} = \infty$ and

$$(4.10) \quad E(\phi_n, \psi_n; D_{8t_n}(x_n) \setminus D_{t_n}(x_n)) \geq \bar{\epsilon} > 0.$$

Setting

$$u_n(x) := \phi_n(x_n + t_n x), v_n(x) := \sqrt{t_n} \psi_n(x_n + t_n x),$$

then it is easy to see that 0 is an energy concentration point for (u_n, v_n) . We have to consider the following two cases:

(a) (u_n, v_n) has no other energy concentration points except 0.

By Theorem 2.1, passing to a subsequence, we may assume that (u_n, v_n) converges to a Dirac-harmonic map $(\sigma, \xi) : \mathbb{R}^2 \rightarrow N$ strongly in $W_{loc}^{1,2}(\mathbb{R}^2) \times L_{loc}^4(\mathbb{R}^2)$ as $n \rightarrow \infty$. In particular, we have

$$\lim_{n \rightarrow \infty} E(u_n, v_n; D_8 \setminus D_1) = E(\sigma, \xi; D_8 \setminus D_1) \geq \bar{\epsilon}.$$

According to the standard theory of Dirac-harmonic maps, we know that (σ, ξ) is a nontrivial Dirac-harmonic sphere. This is the second bubble. This is a contradiction to the ‘‘one bubble’’ assumption.

(b) (u_n, v_n) has another energy concentration point $p \neq 0$.

Without loss of generality, we may assume that p is the only blow-up point in $D_r(p)$ for some small $r > 0$. By the standard theory of blow-up analysis, there exist $x'_n \rightarrow p$ and $r'_n \rightarrow 0$ such that

$$(4.11) \quad E(u_n, v_n; D_{r'_n}(x'_n)) = \sup_{\substack{x \in D_r(p), s \leq r_n \\ D_s(x) \subset D_r(p)}} E(u_n, v_n; D_s(x)) = \frac{\epsilon_0^2}{4}.$$

From the process of constructing the first bubble, we know that there exists a nontrivial Dirac-harmonic sphere $(\tilde{\sigma}, \tilde{\xi})$ such that

$$(u_n(x'_n + r'_n x), r_n^{1/2} v_n(x'_n + r'_n x)) \rightarrow (\tilde{\sigma}, \tilde{\xi}) \text{ strongly in } W_{loc}^{1,2}(\mathbb{R}^2) \times L_{loc}^4(\mathbb{R}^2)$$

as $n \rightarrow \infty$. This is

$$(\phi_n(x_n + t_n x'_n + t_n r'_n x), (t_n r'_n)^{1/2} \psi_n(x_n + t_n x'_n + t_n r'_n x)) \rightarrow (\tilde{\sigma}, \tilde{\xi}) \text{ strongly in } W_{loc}^{1,2}(\mathbb{R}^2) \times L_{loc}^4(\mathbb{R}^2)$$

as $n \rightarrow \infty$. By (4.11), $(\tilde{\sigma}, \tilde{\xi})$ is nontrivial. Therefore, we again get the second bubble contradicting the “one bubble” assumption. So, we proved **Claim** (4.9).

By Theorem 2.1, for any $t \in (r_n R, \delta)$, we obtain

$$\begin{aligned} & \| |x|^{\frac{3}{4}} \nabla \phi_n \|_{L^8(D_{2t}(x_n) \setminus D_t(x_n))} + \| |x|^{\frac{3}{8}} \psi_n \|_{L^{16}(D_{2t}(x_n) \setminus D_t(x_n))} \\ & \leq C(\| \nabla \phi_n \|_{L^2(D_{4t}(x_n) \setminus D_{t/2}(x_n))} + \| \psi_n \|_{L^4(D_{4t}(x_n) \setminus D_{t/2}(x_n))} + t \| \tau_n \|_{L^2(D_{4t}(x_n) \setminus D_{t/2}(x_n))} \\ & \quad + \sqrt{t} \| h_n \|_{L^2(D_{4t}(x_n) \setminus D_{t/2}(x_n))}) \\ (4.12) \quad & \leq C(\epsilon + \sqrt{\delta}). \end{aligned}$$

For simplicity, we will denote $\phi_n, \psi_n, \tau_n, h_n$ by ϕ, ψ, τ and h respectively.

We define $\phi^*(r)$ and $\psi^*(r)$ as follows:

$$(4.13) \quad \phi^*(r) = \frac{1}{2\pi} \int_0^{2\pi} \phi(r, \theta) d\theta \quad \text{and} \quad \psi^*(r) = \frac{1}{2\pi} \int_0^{2\pi} \psi(r, \theta) d\theta.$$

Next, we use the same method as in [17] to compute the equation for $(\phi - \phi^*, \psi - \psi^*)$. Here, for reader's convenience, we repeat this process again.

By equation (1.7), we have

$$\begin{aligned} \Delta \phi^*(t) &= \frac{1}{2\pi} \int_0^{2\pi} -A(\phi)(d\phi, d\phi) + \text{Re}(P(\phi)(\mathcal{A}(\phi)(d\phi(e_\alpha), e_\alpha \cdot \psi); \psi)) + \tau d\theta \\ &= \text{I} + \text{II} + \frac{1}{2\pi} \int_0^{2\pi} \tau d\theta. \end{aligned}$$

Computing directly, we have

$$\begin{aligned} \text{I} &= \frac{1}{2\pi} \int_0^{2\pi} -A(\phi)(d\phi, d\phi) + A(\phi^*)(d\phi, d\phi) - A(\phi^*)(d\phi, d\phi) + A(\phi^*)(d\phi^*, d\phi^*) \\ & \quad - A(\phi^*)(d\phi^*, d\phi^*) d\theta \\ &= -A(\phi^*)(d\phi^*, d\phi^*) + \frac{1}{2\pi} \int_0^{2\pi} A^4(\phi - \phi^*) + A^5 \nabla(\phi - \phi^*) d\theta, \end{aligned}$$

and

$$\begin{aligned}
\Pi &= \frac{1}{2\pi} \operatorname{Re} \int_0^{2\pi} P(\phi) (\mathcal{A}(\phi)(d\phi(e_\alpha), e_\alpha \cdot \psi); \psi) - P(\phi^*) (\mathcal{A}(\phi)(d\phi(e_\alpha), e_\alpha \cdot \psi); \psi) \\
&\quad + P(\phi^*) (\mathcal{A}(\phi)(d\phi(e_\alpha), e_\alpha \cdot \psi); \psi) - P(\phi^*) (\mathcal{A}(\phi^*)(d\phi(e_\alpha), e_\alpha \cdot \psi); \psi) \\
&\quad + P(\phi^*) (\mathcal{A}(\phi^*)(d\phi(e_\alpha), e_\alpha \cdot \psi); \psi) - P(\phi^*) (\mathcal{A}(\phi^*)(d\phi^*(e_\alpha), e_\alpha \cdot \psi); \psi) \\
&\quad + P(\phi^*) (\mathcal{A}(\phi^*)(d\phi^*(e_\alpha), e_\alpha \cdot \psi); \psi) - P(\phi^*) (\mathcal{A}(\phi^*)(d\phi^*(e_\alpha), e_\alpha \cdot \psi^*); \psi) \\
&\quad + P(\phi^*) (\mathcal{A}(\phi^*)(d\phi^*(e_\alpha), e_\alpha \cdot \psi^*); \psi) - P(\phi^*) (\mathcal{A}(\phi^*)(d\phi^*(e_\alpha), e_\alpha \cdot \psi^*); \psi^*) \\
&\quad + P(\phi^*) (\mathcal{A}(\phi^*)(d\phi^*(e_\alpha), e_\alpha \cdot \psi^*); \psi^*) d\theta \\
&= \operatorname{Re}(P(\phi^*) (\mathcal{A}(\phi^*)(d\phi^*(e_\alpha), e_\alpha \cdot \Psi^*); \psi^*)) + \frac{1}{2\pi} \int_0^{2\pi} A^4(\phi - \phi^*) + A^5 \nabla(\phi - \phi^*) \\
&\quad + \frac{1}{2\pi} \operatorname{Re} \int_0^{2\pi} A^6(\psi - \psi^*) d\theta,
\end{aligned}$$

where A^i may differ from line to line and just stands for an expression satisfying

$$|A^4| \leq C(N)(|d\phi|^2 + |d\phi||\psi|^2), \quad |A^5| \leq C(N)(|d\phi| + |\psi|^2), \quad |A^6| \leq C(N)|d\phi||\psi|.$$

Moreover, (4.12) implies

$$\| |x|^{\frac{3}{2}} A^4 \|_{L^4(D_{2t} \setminus D_t)} + \| |x| A^6 \|_{L^4(D_{2t} \setminus D_t)} + \| |x|^{\frac{1}{2}} A^5 \|_{L^4(D_{2t} \setminus D_t)} \leq C\epsilon$$

for any $t \in (\frac{1}{2}r_n R, 2\delta)$. Then, we get

$$\begin{aligned}
\Delta(\phi - \phi^*) &= -A(\phi)(d\phi, d\phi) + A(\phi^*)(d\phi^*, d\phi^*) \\
&\quad + \operatorname{Re} P(\phi) (\mathcal{A}(\phi)(d\phi(e_\alpha), e_\alpha \cdot \psi); \psi) - \operatorname{Re} P(\phi^*) (\mathcal{A}(\phi^*)(d\phi^*(e_\alpha), e_\alpha \cdot \psi^*); \psi^*) \\
&\quad - \frac{1}{2\pi} \int_0^{2\pi} A^4(\phi - \phi^*) + A^5 \nabla(\phi - \phi^*) + \operatorname{Re}(A^6(\psi - \psi^*)) d\theta \\
&\quad + \tau - \frac{1}{2\pi} \int_0^{2\pi} \tau d\theta.
\end{aligned}$$

Using the same method, we get

$$\begin{aligned}
\Delta(\phi - \phi^*) &= A^1(\phi - \phi^*) + A^2 \nabla(\phi - \phi^*) + \operatorname{Re}(A^3(\psi - \psi^*)) \\
&\quad + \frac{1}{2\pi} \int_0^{2\pi} A^4(\phi - \phi^*) + A^5 \nabla(\phi - \phi^*) + \operatorname{Re}(A^6(\psi - \psi^*)) d\theta \\
(4.14) \quad &\quad + \tau - \frac{1}{2\pi} \int_0^{2\pi} \tau d\theta,
\end{aligned}$$

and

$$\begin{aligned}
\delta(\psi - \psi^*) &= B^1(\phi - \phi^*) + B^2 \nabla(\phi - \phi^*) + B^3(\psi - \psi^*) \\
&\quad + \frac{1}{2\pi} \int_0^{2\pi} B^4(\phi - \phi^*) + B^5 \nabla(\phi - \phi^*) + B^6(\psi - \psi^*) d\theta \\
(4.15) \quad &\quad + h - \frac{1}{2\pi} \int_0^{2\pi} h d\theta,
\end{aligned}$$

where $A^i, B^i, i = 1, \dots, 6$ satisfy

$$\begin{aligned} |A^1| + |A^4| &\leq C(N)(|d\phi|^2 + |d\phi||\psi|^2), \quad |A^2| + |A^5| + |B^3| + |B|^6 \leq C(N)(|d\phi| + |\psi|^2), \\ |A^3| + |A^6| + |B^1| + |B^4| &\leq C(N)|d\phi||\psi|, \quad |B^2| + |B^5| \leq C(N)|\psi|. \end{aligned}$$

and

$$(4.16) \quad \begin{aligned} &\| |x|^{\frac{3}{2}}(|A^1| + |A^4|) \|_{L^4(D_{2r} \setminus D_r)} + \| |x|(|A^3| + |A^6| + |B^1| + |B^4|) \|_{L^4(D_{2r} \setminus D_r)} \\ &+ \| |x|^{\frac{1}{2}}(|A^2| + |A^5| + |B^3| + |B^6|) \|_{L^4(D_{2r} \setminus D_r)} + \| |B^2| + |B^5| \|_{L^4(D_{2r} \setminus D_r)} \leq C\epsilon \end{aligned}$$

for any $t \in (\frac{1}{2}r_n R, 2\delta)$ by (4.12).

Without loss of generality, we may assume $\delta = e^{m_n L} r_n R$ for some positive integer m_n which tends to ∞ as $n \rightarrow \infty$. Substituting $u = \phi - \phi^*$ and $v = \psi - \psi^*$ in Corollary 3.5, we obtain the energy decay in the θ -direction,

$$(4.17) \quad \begin{aligned} \| r^{-1} \frac{\partial \phi}{\partial \theta} \|_{L^2(P_i)} + \| r^{-1} \frac{\partial \psi}{\partial \theta} \|_{L^{\frac{4}{3}}(P_i)} &\leq \| \nabla u \|_{L^2(P_i)} + \| \nabla v \|_{L^{\frac{4}{3}}(P_i)} \\ &\leq C((e^{iL} r_n R)^{\frac{1}{2}} + F_0^{1/2}(u, v) + F_{m_n}^{1/2}(u, v))(e^{-\frac{1}{2}(m_n - i)L} + e^{-\frac{1}{2}iL}) \\ &\leq C(\sqrt{\epsilon} + \sqrt{\delta})(e^{-\frac{1}{2}(m_n - i)L} + e^{-\frac{1}{2}iL}), \end{aligned}$$

where the last inequality follows from Poincaré's inequality and the assumption (4.9).

By Corollary 2.4, we get

$$(4.18) \quad \begin{aligned} \| \frac{\partial \phi}{\partial r} \|_{L^2(P_i)} &\leq \| r^{-1} \frac{\partial \phi}{\partial \theta} \|_{L^2(P_i)}^{\frac{1}{2}} + C \| r^{-1} \frac{\partial \psi}{\partial \theta} \|_{L^{\frac{4}{3}}(P_i)}^{\frac{1}{2}} + C \sqrt{e^{iL} r_n R} \\ &\leq C(\epsilon^{\frac{1}{4}} + \delta^{\frac{1}{4}})(e^{-\frac{1}{4}(m_n - i)L} + e^{-\frac{1}{4}iL}). \end{aligned}$$

Therefore,

$$(4.19) \quad \| \nabla \phi \|_{L^2(P_i)} \leq C(\epsilon^{\frac{1}{4}} + \delta^{\frac{1}{4}})(e^{-\frac{1}{4}(m_n - i)L} + e^{-\frac{1}{4}iL}).$$

Then, by Theorem 2.1, we have

$$(4.20) \quad E(\phi; D_\delta(x_n) \setminus D_{r_n R}(x_n)) \leq \sum_{i=0}^{m_n - 1} \| \nabla \phi \|_{L^2(P_i)}^2 \leq C(\epsilon^{\frac{1}{2}} + \delta^{\frac{1}{2}})$$

and

$$(4.21) \quad OSC_{D_\delta(x_n) \setminus D_{r_n R}(x_n)} \phi_n \leq C \sum_{i=0}^{m_n - 1} (\| \nabla \phi \|_{L^2(P_i)} + e^{iL} r_n R \| \tau \|_{L^2(P_i)}) \leq C(\epsilon^{\frac{1}{4}} + \delta^{\frac{1}{4}}).$$

So, we have proved (4.1) and (4.3).

Combining this with equation (1.8), we get

$$\begin{aligned} \| \frac{\partial \psi}{\partial r} \|_{L^{\frac{4}{3}}(P_i)} &\leq \| r^{-1} \frac{\partial \psi}{\partial \theta} \|_{L^{\frac{4}{3}}(P_i)} + C \| \nabla \phi \|_{L^2(P_i)} \| \psi \|_{L^4(P_i)} + \| h \|_{L^{\frac{4}{3}}(P_i)} \\ &\leq C(\epsilon^{\frac{1}{4}} + \delta^{\frac{1}{4}})(e^{-\frac{1}{4}(m_n - i)L} + e^{-\frac{1}{4}iL}) + C e^{iL} r_n R \| h \|_{L^4(P_i)} \\ &\leq C(\epsilon^{\frac{1}{4}} + \delta^{\frac{1}{4}})(e^{-\frac{1}{4}(m_n - i)L} + e^{-\frac{1}{4}iL}). \end{aligned}$$

Thus,

$$(4.22) \quad \|\nabla\psi\|_{L^{\frac{4}{3}}(P_i)} \leq C(\epsilon^{\frac{1}{4}} + \delta^{\frac{1}{4}})(e^{-\frac{1}{4}(m_n-i)L} + e^{-\frac{1}{4}iL}).$$

Taking a cut-off function $\eta \in C_0^\infty(D_\delta)$, such that $0 \leq \eta \leq 1$, $\eta \equiv 1$ in $D_{\frac{1}{2}\delta} \setminus D_{2r_n R}$ and

$$|\nabla\eta| \leq \frac{C}{\delta} \text{ in } D_\delta \setminus D_{\frac{1}{2}\delta} \text{ and } |\nabla\eta| \leq \frac{C}{r_n R} \text{ in } D_{2r_n R} \setminus D_{r_n R},$$

by the elliptic estimates for first order equations and Sobolev embedding, we obtain

$$\begin{aligned} \|\eta\psi\|_{L^4(D_\delta)} &\leq C\|\not\partial(\eta\psi)\|_{L^{\frac{4}{3}}(D_\delta)} \\ &\leq C\left(\frac{1}{\delta}\|\psi\|_{L^{\frac{4}{3}}(D_\delta \setminus D_{\frac{1}{2}\delta})} + \frac{1}{r_n R}\|\psi\|_{L^{\frac{4}{3}}(D_{2r_n R} \setminus D_{r_n R})} + \|\nabla\psi\|_{L^{\frac{4}{3}}(D_\delta \setminus D_{r_n R})} + \|h\|_{L^{\frac{4}{3}}(D_\delta)}\right) \\ &\leq C\left(\|\psi\|_{L^4(D_\delta \setminus D_{\frac{1}{2}\delta})} + \|\psi\|_{L^4(D_{2r_n R} \setminus D_{r_n R})} + \sum_{i=1}^{m_n} \|\nabla\psi\|_{L^{\frac{4}{3}}(P_i)} + \delta\|h\|_{L^4(D_\delta)}\right) \\ &\leq C(\epsilon^{\frac{1}{4}} + \delta^{\frac{1}{4}}), \end{aligned}$$

where the last inequality follows from (4.22). This is

$$(4.23) \quad E(\psi; D_\delta(x_n) \setminus D_{r_n R}(x_n)) \leq C(\epsilon + \delta).$$

This is (4.2) and we finished the proof of Theorem 4.1. \square

Proof of Theorem 1.2: It is easy to see that Theorem 1.2 is a consequence of Theorem 4.1, the removable singularity Theorem 2.5 and the standard argument in [7]. \square

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